



# DISSIPATIVE DISTRIBUTED SYSTEMS

Curtainday, October 19, 2001

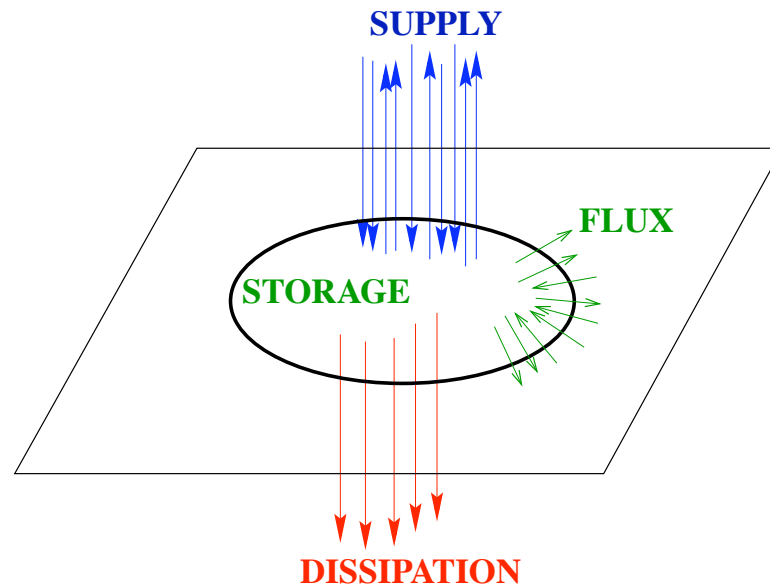
RUG

A dissipative system absorbs ‘supply’, ‘globally’ over time and space.

∴ Can this be expressed ‘locally’, as

**rate of change in storage + spatial flux  $\leq$  supply rate**

**= supply rate + dissipation rate ??**



$\Rightarrow$ , for 1-D state space systems, the Kalman-Yacubovich-Popov lemma.



$\Rightarrow$  Lyapunov theory for **open** dynamical systems, **LMI's**, their many applications.

$\Rightarrow$  applications in the analysis of open physical systems, synthesis procedures, robust control, passivation control.

**Joint work with Harish Pillai (IIT, Mumbai)**



**& Harry Trentelman.**



# **BEHAVIORAL SYSTEMS**

A system :=  $\Sigma = (\mathbb{T}, \mathbb{W}, \mathcal{B})$

$\mathbb{T}$  = the set of independent variables  
time, space, time and space

$\mathbb{W}$  = the set of dependent variables  
(= space where the variables take on their values,  
signal space, space of field variables, . . .)

$\mathcal{B} \subseteq \mathbb{W}^{\mathbb{T}}$  : the behavior = the admissible trajectories

$$\Sigma = (\mathbb{T}, \mathbb{W}, \mathfrak{B})$$

for a trajectory  $w : \mathbb{T} \rightarrow \mathbb{W}$ , we thus have:

$w \in \mathfrak{B}$  : the model **allows** the trajectory  $w$ ,

$w \notin \mathfrak{B}$  : the model **forbids** the trajectory  $w$ .

In this lecture,  $\mathbb{T} = \mathbb{R}^n$ , (**'n-D systems'**),  $\mathbb{W} = \mathbb{R}^w$ ,

$w : \mathbb{R}^n \rightarrow \mathbb{R}^w, (w_1(x_1, \dots, x_n), \dots, w_w(x_1, \dots, x_n))$ ,

often,  $n = 4$ , independent variables  $(t, x, y, z)$ ,

$\mathfrak{B}$  = solutions of a system of constant coefficient  
linear PDE's.

**'Linear (shift-invariant distributed) differential systems'**.

**Example: *Maxwell's equations***



$$\begin{aligned}\nabla \cdot \vec{E} &= \frac{1}{\epsilon_0} \rho, \\ \nabla \times \vec{E} &= -\frac{\partial}{\partial t} \vec{B}, \\ \nabla \cdot \vec{B} &= 0, \\ c^2 \nabla \times \vec{B} &= \frac{1}{\epsilon_0} \vec{j} + \frac{\partial}{\partial t} \vec{E}.\end{aligned}$$

$\mathbb{T} = \mathbb{R} \times \mathbb{R}^3$  (time and space),

$w = (\vec{E}, \vec{B}, \vec{j}, \rho)$

(electric field, magnetic field, current density, charge density),

$\mathbb{W} = \mathbb{R}^3 \times \mathbb{R}^3 \times \mathbb{R}^3 \times \mathbb{R}$ ,

$\mathfrak{B} =$  set of solutions to these PDE's.

**Note: 10 variables, 8 equations!  $\Rightarrow \exists$  free variables.**



## n-D LINEAR DIFFERENTIAL SYSTEMS

$$T = \mathbb{R}^n, \quad W = \mathbb{R}^w,$$

$\mathfrak{B}$  = **the solutions of a linear constant coefficient system of PDE's.**

Let  $R \in \mathbb{R}^{\bullet \times w}[\xi_1, \dots, \xi_n]$ , and consider

$$R\left(\frac{\partial}{\partial x_1}, \dots, \frac{\partial}{\partial x_n}\right)w = 0. \quad (*)$$

Define the associated behavior

$$\mathfrak{B} = \{w \in \mathcal{C}^\infty(\mathbb{R}^n, \mathbb{R}^w) \mid (*) \text{ holds} \}$$

$\mathcal{C}^\infty(\mathbb{R}^n, \mathbb{R}^w)$  **mainly** for convenience, but important for some results.

Examples: Maxwell's eq'ns, diffusion eq'n, wave eq'n, . . .

**Notation for n-D linear differential systems:**

$$(\mathbb{R}^n, \mathbb{R}^w, \mathfrak{B}) \in \mathfrak{L}_n^w, \quad \text{or } \mathfrak{B} \in \mathfrak{L}_n^w,$$

Note       $\mathfrak{B} = \ker(R(\frac{\partial}{\partial x_1}, \dots, \frac{\partial}{\partial x_n})),$

$$R(\frac{\partial}{\partial x_1}, \dots, \frac{\partial}{\partial x_n})w = 0$$

is called a **'kernel representation'** of  $\mathfrak{B} \in \mathfrak{L}_n^w$ .

THE STRUCTURE OF  $\mathcal{L}_n^w$

## PROPERTIES of $\mathcal{L}_n^W$

- $\exists 1 \leftrightarrow 1$  relation between  $\mathcal{L}_n^W$  and the **submodules** of  $\mathbb{R}^W[\xi_1, \dots, \xi_n]$
- **Elimination theorem**: the behavior of a subset of the system variables is also described by a PDE
- **Controllability**; image representations
- **Observability**

Work of Shankar/Pillai, Oberst.



**DESCRIBE  $(\vec{E}, \vec{j})$  IN MAXWELL'S EQ'NS ?**

**Eliminate  $\vec{B}, \rho$  from Maxwell's equations. Straightforward computation yields**

$$\begin{aligned}\epsilon_0 \frac{\partial}{\partial t} \nabla \cdot \vec{E} + \nabla \cdot \vec{j} &= 0, \\ \epsilon_0 \frac{\partial^2}{\partial t^2} \vec{E} + \epsilon_0 c^2 \nabla \times \nabla \times \vec{E} + \frac{\partial}{\partial t} \vec{j} &= 0.\end{aligned}$$

**Elimination theorem  $\Rightarrow$  this exercise would be exact & successful.**

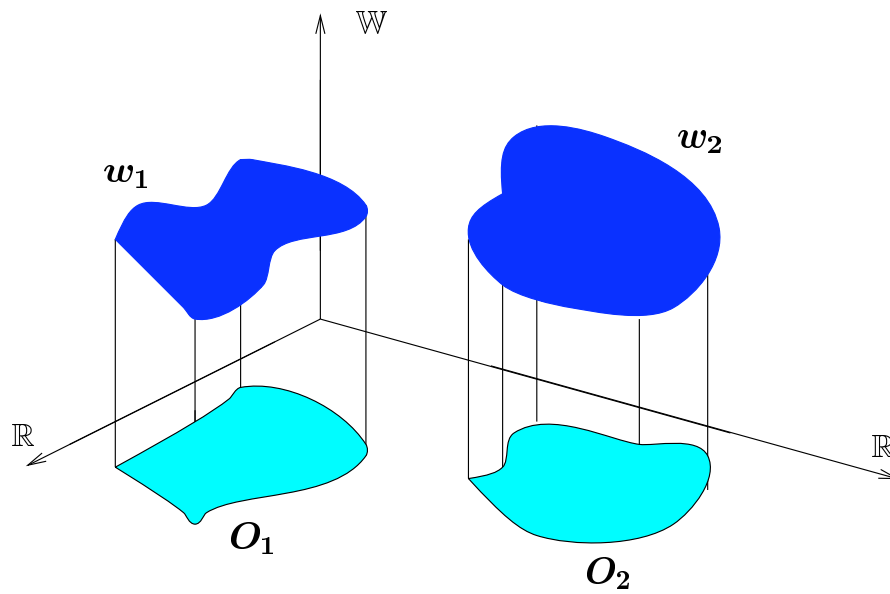
## CONTROLLABILITY

**Definition:**  $\mathfrak{B} \in \mathcal{L}_n^w$  is said to be

*controllable*

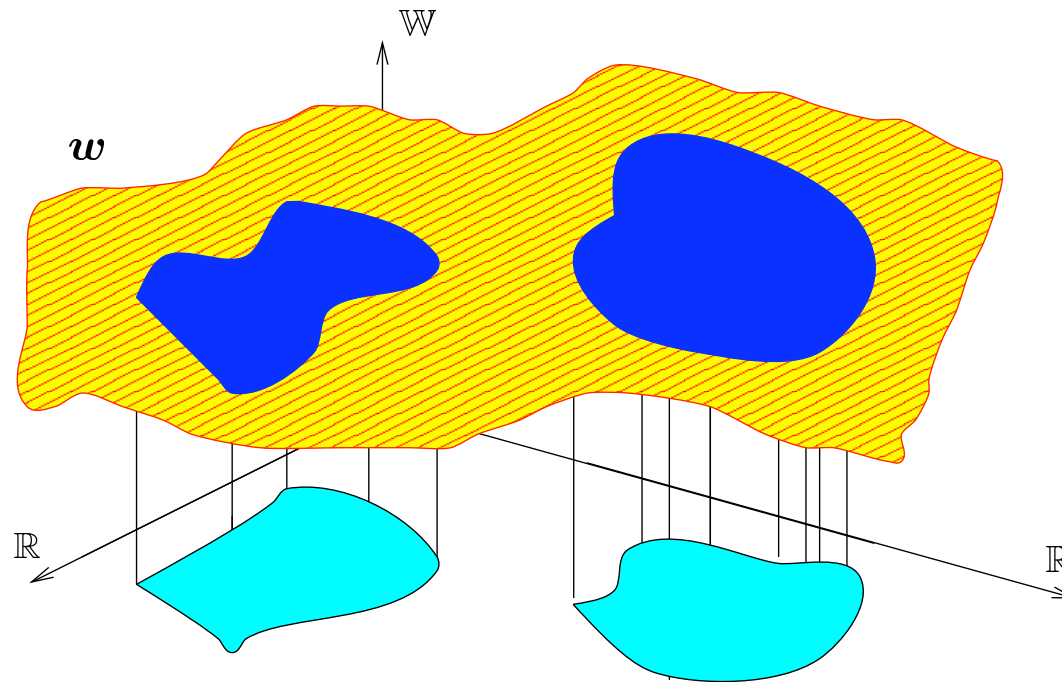
if ...

In pictures:



$w_1, w_2 \in \mathcal{B}.$

$w \in \mathcal{B}$  ‘patches’  $w_1, w_2 \in \mathcal{B}$ .



**Controllability**  $:\Leftrightarrow$  ‘patch-ability’.

Special case: **Kalman controllability** for state space systems.



## CONDITIONS FOR CONTROLLABILITY

$$R\left(\frac{\partial}{\partial x_1}, \dots, \frac{\partial}{\partial x_n}\right)w = 0,$$

is called a *kernel representation* of  $\mathfrak{B} = \ker\left(R\left(\frac{\partial}{\partial x_1}, \dots, \frac{\partial}{\partial x_n}\right)\right)$ ;

$$w = M\left(\frac{\partial}{\partial x_1}, \dots, \frac{\partial}{\partial x_n}\right)\ell,$$

is called an *image representation* of  $\mathfrak{B} = \text{im}\left(M\left(\frac{\partial}{\partial x_1}, \dots, \frac{\partial}{\partial x_n}\right)\right)$ .

**Elimination theorem**  $\Rightarrow$  every image is also a **kernel**.

?? Which kernels are also images ??

Theorem (Pillai/Shankar):

The following are equivalent for  $\mathfrak{B} \in \mathcal{L}_n^w$  :

1.  $\mathfrak{B}$  is **controllable**,
2.  $\mathfrak{B}$  admits an image representation,

⋮

## ARE MAXWELL'S EQUATIONS CONTROLLABLE ?

The following well-known equations in the *scalar potential*  $\phi : \mathbb{R} \times \mathbb{R}^3 \rightarrow \mathbb{R}$  and the *vector potential*  $\vec{A} : \mathbb{R} \times \mathbb{R}^3 \rightarrow \mathbb{R}^3$ , generate exactly the solutions to Maxwell's equations:

$$\vec{E} = -\frac{\partial}{\partial t} \vec{A} - \nabla \phi,$$

$$\vec{B} = \nabla \times \vec{A},$$

$$\vec{j} = \epsilon_0 \frac{\partial^2}{\partial t^2} \vec{A} - \epsilon_0 c^2 \nabla^2 \vec{A} + \epsilon_0 c^2 \nabla (\nabla \cdot \vec{A}) + \epsilon_0 \frac{\partial}{\partial t} \nabla \phi,$$

$$\rho = -\epsilon_0 \frac{\partial}{\partial t} \nabla \cdot \vec{A} - \epsilon_0 \nabla^2 \phi.$$

Proves controllability. Illustrates the interesting connection

**controllability  $\Leftrightarrow \exists$  a potential!**

## OBSERVABILITY

$w = M\left(\frac{\partial}{\partial x_1}, \dots, \frac{\partial}{\partial x_n}\right)\ell$  is said to be **observable** if  $\ell$  can be deduced from  $w$ , i.e., if  $M\left(\frac{\partial}{\partial x_1}, \dots, \frac{\partial}{\partial x_n}\right)$  is **injective**.

∴ Controllability  $\Rightarrow \exists$  an observable image representation ??

For 1-D systems, **yes!**

For n-D systems, **not necessarily!**

**Non-example:** Maxwell's equations. Potential is not observable!

**No potential ever is, for Maxwell's equations.**

For n-D systems: image representation requires **'hidden'** variables.

# DISSIPATIVITY

**Multi-index notation:**

$$\mathbf{x} = (x_1, \dots, x_n),$$

$$\mathbf{k} = (k_1, \dots, k_n), \ell = (\ell_1, \dots, \ell_n),$$

$$\boldsymbol{\xi} = (\xi_1, \dots, \xi_n), \boldsymbol{\zeta} = (\zeta_1, \dots, \zeta_n), \boldsymbol{\eta} = (\eta_1, \dots, \eta_n),$$

$$\frac{d}{dx} = \left( \frac{\partial}{\partial x_1}, \dots, \frac{\partial}{\partial x_n} \right), \frac{d^{\mathbf{k}}}{dx^{\mathbf{k}}} = \left( \frac{\partial^{k_1}}{\partial x_1^{k_1}}, \dots, \frac{\partial^{k_n}}{\partial x_n^{k_n}} \right),$$

$$dx = dx_1 dx_2 \dots dx_n,$$

$$R\left(\frac{d}{dx}\right)w = 0 \quad \text{for} \quad R\left(\frac{\partial}{\partial x_1}, \dots, \frac{\partial}{\partial x_n}\right)w = 0,$$

etc.

## QDF's

The quadratic map in the variables  $w$  and their partial derivatives, defined by

$$w \mapsto \sum_{k,l} \left( \frac{d^k}{dx^k} w \right)^\top \Phi_{k,l} \left( \frac{d^l}{dx^l} w \right)$$

is called quadratic differential form (QDF) on  $\mathcal{C}^\infty(\mathbb{R}^n, \mathbb{R}^w)$ .

Here  $\Phi_{k,l} \in \mathbb{R}^{w \times w}$

(given matrices, only finite number  $\neq 0$ ,  $\Phi_{k,l} = \Phi_{l,k}^\top$ ).

**Introduce the  $2n$ -variable polynomial matrix  $\Phi$  defined by**

$$\Phi(\zeta, \eta) = \sum_{k, \ell} \Phi_{k, \ell} \zeta^k \eta^\ell.$$

**Denote this QDF as  $Q_\Phi$ ; whence**

$$Q_\Phi : w \in \mathcal{C}^\infty(\mathbb{R}^n, \mathbb{R}^w) \mapsto Q_\Phi(w) \in \mathcal{C}^\infty(\mathbb{R}^n, \mathbb{R}).$$

**A QDF is parameterized by a  $\Phi(\zeta, \eta) \in \mathbb{R}^{w \times w}[\zeta, \eta]$ .**



## DISSIPATIVE DISTRIBUTED SYSTEMS

We consider only **controllable linear differential systems** and **QDF's**.

**Definition:**  $\mathfrak{B} \in \mathcal{L}_n^w$ , controllable, is said to be **dissipative** with respect to the **supply rate**  $Q_\Phi$  (a QDF) if

$$\int_{\mathbb{R}^n} Q_\Phi(w) dx \geq 0$$

for all  $w \in \mathfrak{B}$  of compact support, i.e., for all  $w \in \mathfrak{B} \cap \mathcal{D}$ .

**Idea:**  $Q_{\Phi}(w)(x_1, \dots, x_n) dx_1 \cdots dx_n$  :  
rate of 'energy' delivered to the system.

**Dissipativity** :  $\Leftrightarrow$

$$\int_{\mathbb{R}} \left( \int_{\mathbb{R}^3} Q_{\Phi}(w) dx dy dz \right) dt \geq 0 \quad \text{for all } w \in \mathfrak{B}.$$

A dissipative system **absorbs** net energy (integrated over space and time).

## Examples:

### Maxwell's eq'ns:

dissipative (in fact, conservative) w.r.t. the QDF  $-\vec{E} \cdot \vec{j}$

### Passive electrical circuits

### Mechanical systems

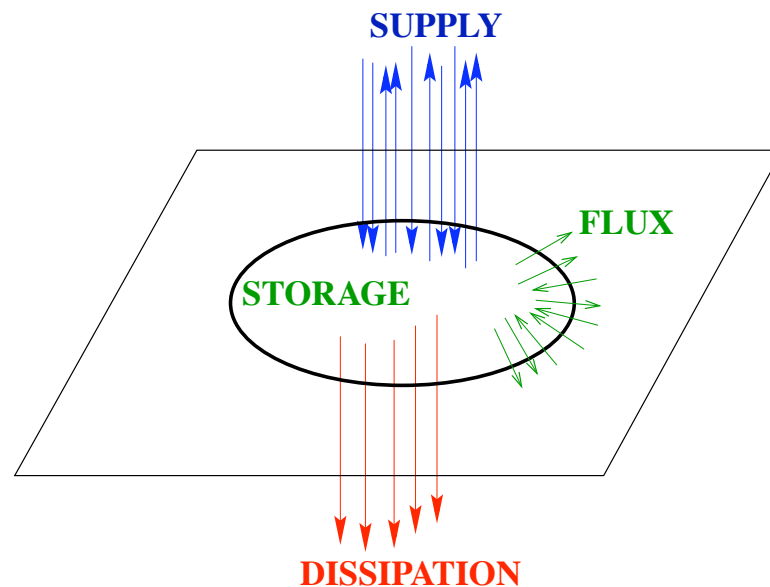
### Thermodynamic systems

## LOCAL VERSION

Assume that a system is 'globally' dissipative.

∴ Can this dissipativity be expressed through a 'local' law??

$$\frac{d}{dt} \text{Storage} + \text{Spatial flux} \leq \text{Supply.}$$



**Supply = Stored + radiated + dissipated**

$$\int_{\mathbb{R}^n} Q_{\Phi}(w) \geq 0, \quad \text{for all } w \in \mathfrak{B} \cap \mathfrak{D} \quad \text{‘global dissipativity’}$$

? ⇕ ?

$$\frac{\partial}{\partial t} Q_{\Phi_t}(w) + \frac{\partial}{\partial x} Q_{\Phi_x}(w) + \frac{\partial}{\partial y} Q_{\Phi_y}(w) + \frac{\partial}{\partial z} Q_{\Phi_z}(w) \leq Q_{\Phi}(w)$$

for all  $w \in \mathfrak{B}$  ‘local dissipativity’.

⇐: easy

⇒: !! construct  $\Phi_t, \Phi_x, \Phi_y, \Phi_z$  from  $\mathfrak{B}$  and  $\Phi$  !!

## **THE LOCAL DISSIPATION LAW**

Main Theorem: Let  $\mathfrak{B} \in \mathcal{L}_n^w$  be controllable.

Then  $\mathfrak{B}$  is **dissipative** with respect to the *supply rate*  $Q_\Phi$   
iff

there exist

an **image representation**  $w = M\left(\frac{d}{dx}\right)\ell$  of  $\mathfrak{B}$ , and

an **n–vector of QDF’s**  $Q_\Psi = (Q_{\Psi_1}, \dots, Q_{\Psi_n})$

on  $\mathcal{C}^\infty(\mathbb{R}^n, \mathbb{R}^{\dim(\ell)})$ , called the *flux*,

such that the *local dissipation law*

$$\nabla \cdot Q_\Psi(\ell) \leq Q_\Phi(w)$$

holds for all  $(w, \ell)$  that satisfy  $w = M\left(\frac{d}{dx}\right)\ell$ .

As usual  $\nabla \cdot Q_\Psi := \frac{\partial}{\partial x_1} Q_{\Psi_1} + \dots + \frac{\partial}{\partial x_n} Q_{\Psi_n}$ .

Note: the local law involves

(possibly unobservable, - i.e., **hidden!**) latent variables (the  $\ell$ 's).

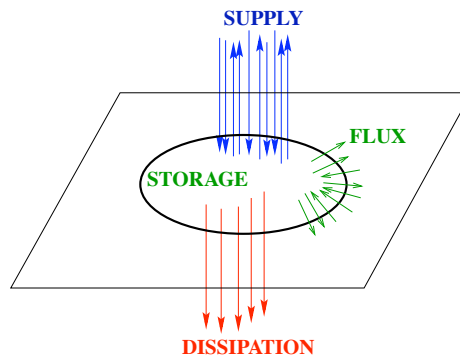


When the variables are  $(t, x, y, z)$ , can be reformulated as:

$\exists$  an image representation  $w = M\left(\frac{\partial}{\partial t}, \frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z}\right)\ell$  of  $\mathfrak{B}$ ,  
a QDF  $S$ , the *storage*, and  
a vector of QDF's,  $F = (F_x, F_y, F_z)$ , the *spatial flux*,  
such that

$$\frac{\partial}{\partial t}S(\ell) + \frac{\partial}{\partial x}F_x(\ell) + \frac{\partial}{\partial y}F_y(\ell) + \frac{\partial}{\partial z}F_z(\ell) \leq Q_{\Phi}(w)$$

holds for all  $(w, \ell)$  that satisfy  $w = M\left(\frac{\partial}{\partial t}, \frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z}\right)\ell$ .



## EXAMPLE: ENERGY STORED IN EM FIELDS

Maxwell's equations are dissipative (in fact, conservative) with respect to  $-\vec{E} \cdot \vec{j}$ , the rate of energy supplied.

Introduce the *stored energy density*,  $S$ , and

the *energy flux density (the Poynting vector)*,  $\vec{F}$ ,

$$S(\vec{E}, \vec{B}) = \frac{\epsilon_0}{2} \vec{E} \cdot \vec{E} + \frac{\epsilon_0 c^2}{2} \vec{B} \cdot \vec{B},$$

$$\vec{F}(\vec{E}, \vec{B}) = \epsilon_0 c^2 \vec{E} \times \vec{B}.$$

The following is a local conservation law for Maxwell's equations:

$$\frac{\partial}{\partial t} S(\vec{E}, \vec{B}) + \nabla \cdot \vec{F}(\vec{E}, \vec{B}) = -\vec{E} \cdot \vec{j}.$$

Local version involves  $\vec{B}$ , unobservable from  $\vec{E}$  and  $\vec{j}$ , the variables in the rate of energy supplied.

## IDEA of the PROOF

Consider  $\mathfrak{B} \in \mathfrak{L}_n^w$ , **controllable**, and  $\Phi \in \mathbb{R}^{w \times w}[\zeta, \eta]$ .

$\mathfrak{B}$  allows an image representation  $w = \mathcal{M}(dx)\ell$ .

Define  $\Phi'(\zeta, \eta) = M^\top(\zeta)\Phi(\zeta, \eta)(M(\eta))$ . Note that

$$Q_\Phi(w) = Q_{\Phi'}(\ell).$$

So we may as well study

$$\int_{-\infty}^{\infty} Q_{\Phi'}(\ell) dt \quad \text{for } \ell \in \mathfrak{D}$$

instead of

$$\int_{-\infty}^{\infty} Q_\Phi(w) dt \quad \text{for } w \in \mathfrak{B} \cap \mathfrak{D}$$

$\rightsquigarrow$  work with  $\Phi'$ , and 'free' signals  $\ell$ .

WLOG:  $\mathfrak{B} = \mathfrak{C}^\infty(\mathbb{R}^n, \mathbb{R}^w) \dots$

$$\int_{\mathbb{R}^n} Q_{\Phi}(w) \geq 0 \text{ for all } w \in \mathfrak{D}$$

$\Updownarrow$  (Parseval)

$$\Phi(-i\omega, i\omega) \geq 0 \text{ for all } \omega \in \mathbb{R}^n$$

$\Updownarrow$  **(Factorization equation)**

$$\exists D : \Phi(-\xi, \xi) = D^{\top}(-\xi)D(\xi)$$

$\Updownarrow$  (easy)

$$\exists \Psi, D : \Psi(\zeta, \eta) = \frac{\Phi(\zeta, \eta) - D^{\top}(\zeta)D(\eta)}{\zeta + \eta}$$

$\Updownarrow$  (clearly)

$$\exists \Psi : Q_{\Psi}(w) \leq Q_{\Phi}(w) \text{ for all } w \in \mathfrak{C}^{\infty}$$

## THE FACTORIZATION EQUATION

Consider

$$X^T(-\xi)X(\xi) = Y(\xi)$$

with  $Y \in \mathbb{R}^{\bullet \times \bullet}[\xi]$  given, and  $X$  the unknown. Solvable??

$\mathbb{R}$

$$X^T(\xi)X(\xi) = Y(\xi)$$

with  $Y \in \mathbb{R}^{\bullet \times \bullet}[\xi]$  given, and  $X$  the unknown.

Under what conditions on  $Y$  does there exist a solution  $X$ ?

Scalar case: !! write  $Y$  as a sum of squares

$$Y = x_1^2 + x_2^2 + \cdots + x_k^2.$$

$$X^T(\xi)X(\xi) = Y(\xi)$$

For  $n = 1$  and  $Y \in \mathbb{R}[\xi]$ , solvable (for  $X \in \mathbb{R}^2[\xi]$  !) iff

$$Y(\alpha) \geq 0 \quad \text{for all } \alpha \in \mathbb{R}.$$

For  $n = 1$ , and  $Y \in \mathbb{R}^{\bullet \times \bullet}[\xi]$ , it is well-known (but non-trivial) that this factorization equation is solvable (with  $X \in \mathbb{R}^{\bullet \times \bullet}[\xi]$  !) iff

$$Y(\alpha) = Y^T(\alpha) \geq 0 \quad \text{for all } \alpha \in \mathbb{R}.$$

For  $n > 1$ , and under these obvious positivity requirements, this equation **can nevertheless** in general not be solved over the **polynomial matrices**, for  $X \in \mathbb{R}^{\bullet \times \bullet}[\xi]$ , but it can be solved over the **matrices of rational functions**, i.e., for  $X \in \mathbb{R}^{\bullet \times \bullet}(\xi)$ .

This factorizability is a consequence of **Hilbert's 17-th problem!**



Solve  $p = p_1^2 + p_2^2 + \cdots + p_k^2$ ,  $p$  given

A polynomial  $p \in \mathbb{R}[\xi_1, \dots, \xi_n]$ , with  $p(\alpha_1, \dots, \alpha_n) \geq 0$  for all  $(\alpha_1, \dots, \alpha_n) \in \mathbb{R}^n$  can in general **not** be expressed as a sum of squares with the  $p_i$ 's  $\in \mathbb{R}[\xi_1, \dots, \xi_n]$  (it can for  $n = 1$ ).

But a rational function (and hence a polynomial)

$p \in \mathbb{R}(\xi_1, \dots, \xi_n)$ , with  $p(\alpha_1, \dots, \alpha_n) \geq 0$ , for all  $(\alpha_1, \dots, \alpha_n) \in \mathbb{R}^n$ , **can** be expressed as a sum of squares, with the  $p_i$ 's  $\in \mathbb{R}(\xi_1, \dots, \xi_n)$ , in fact, with  $k = 2^n$ .



**This solvability readily leads to solvability of the factorization equation over  $\mathbb{R}(\xi_1, \dots, \xi_n)$ , for any  $n$ .**

**The need to introduce rational functions (together with the image representation) are the cause of the **unavoidable** presence of the (possibly unobservable, i.e., **'hidden'**) latent variables in the local dissipation law.**

## CONCLUSIONS

- $\exists$  a nice, **algebraic**, theory of linear shift-invariant differential systems (PDE's)
- **Controllability**  $\Leftrightarrow \exists$  image representation
- **global dissipation**  $\Leftrightarrow \exists$  **local dissipation law**
- Involves **hidden** latent variables (e.g.  $\vec{B}$  applied to Maxwell's eq'ns)
- The proof  $\cong$  **Hilbert's 17-th problem**

More info, ms, copy sheets? Surf to

<http://www.math.rug.nl/~willems>



**HAPPY BIRTHDAY, RUTH !**